

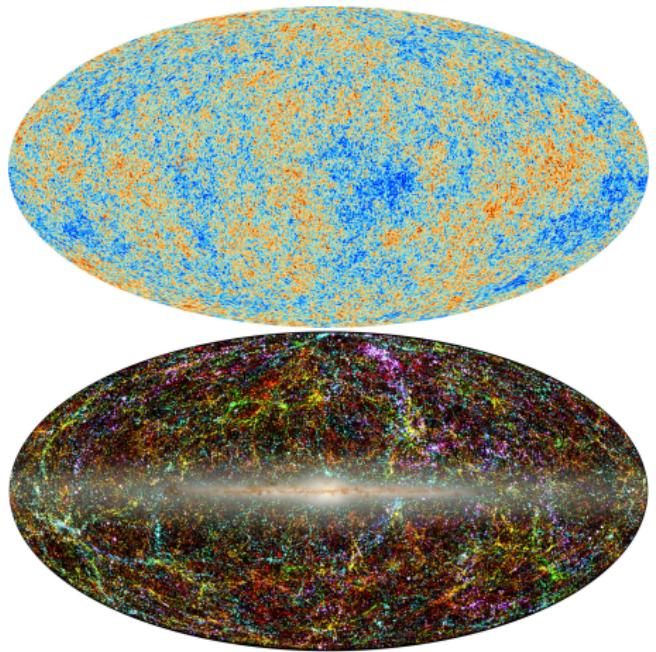
Cosmic Structure Formation With Analytic Methods

Matthias Bartelmann

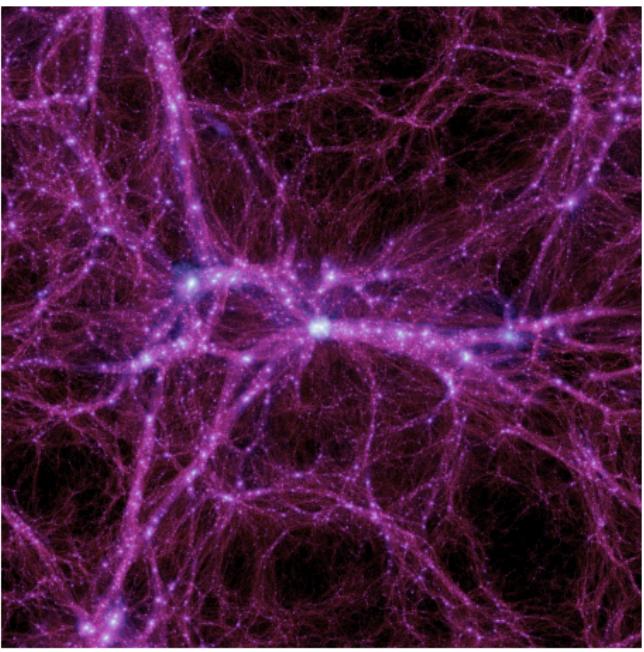
Institute for Theoretical Physics, U. Heidelberg

(Virtual) Seminar, Department of Astronomy, Tsinghua University
Nov. 19th, 2020

Cosmic structure formation

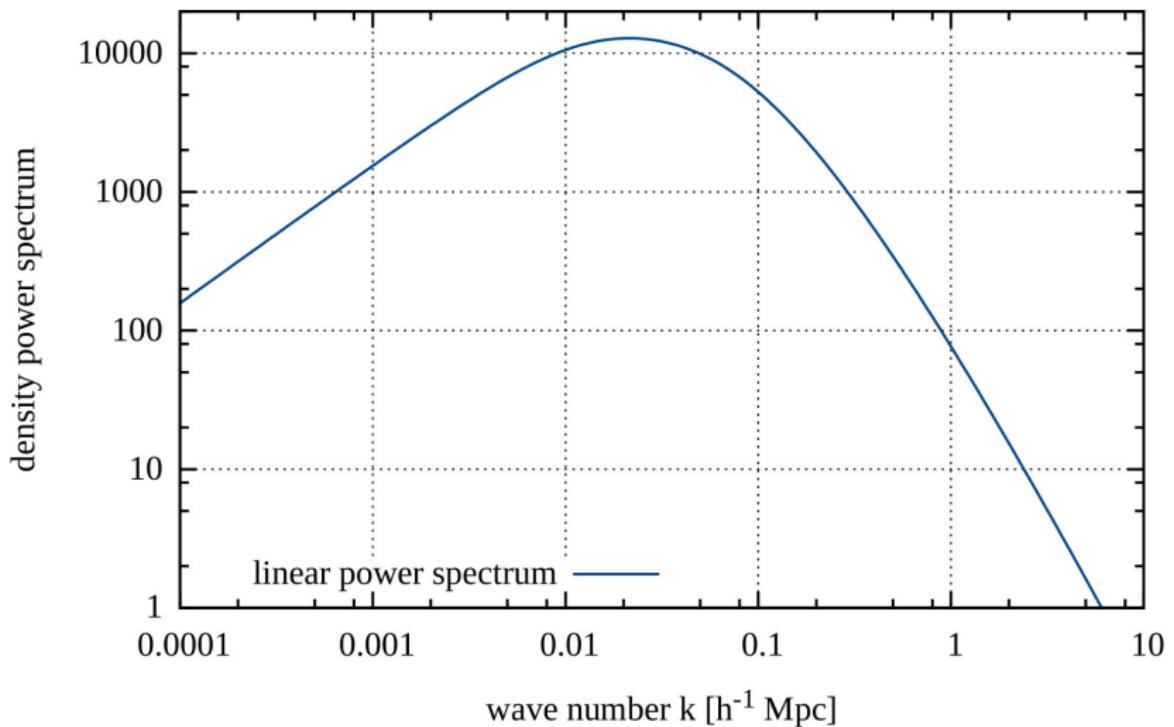


Planck, 2-MASS

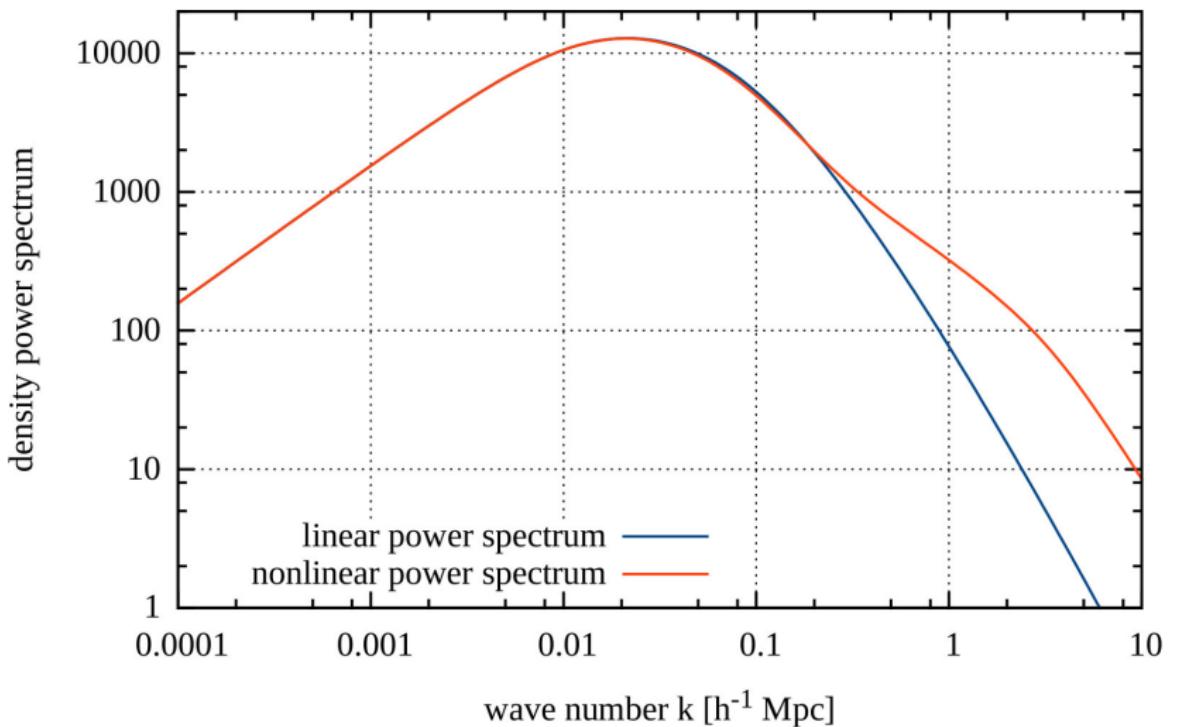


Boylan-Kolchin et al.

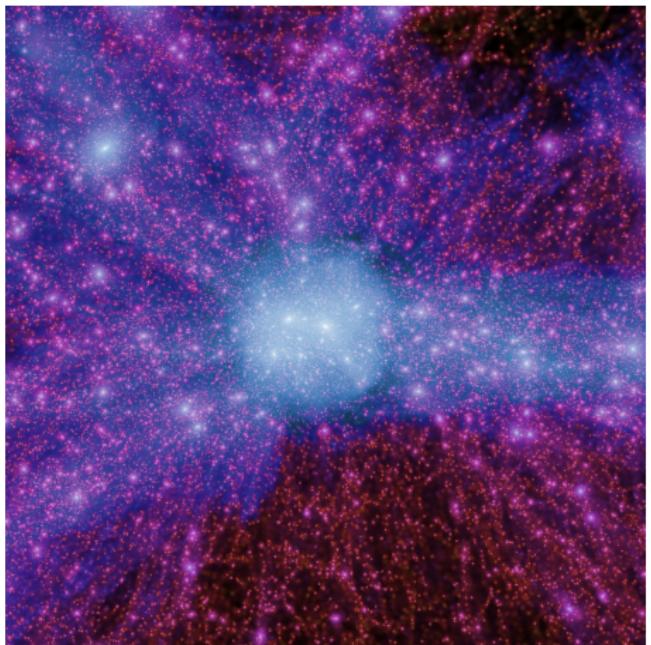
Problems in cosmic structure formation



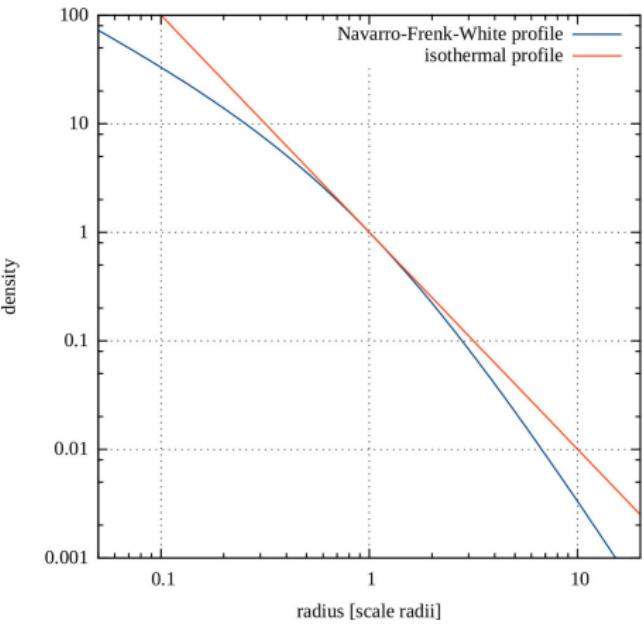
Problems in cosmic structure formation



Problems in cosmic structure formation



Boylan-Kolchin et al.



Problems in cosmic structure formation

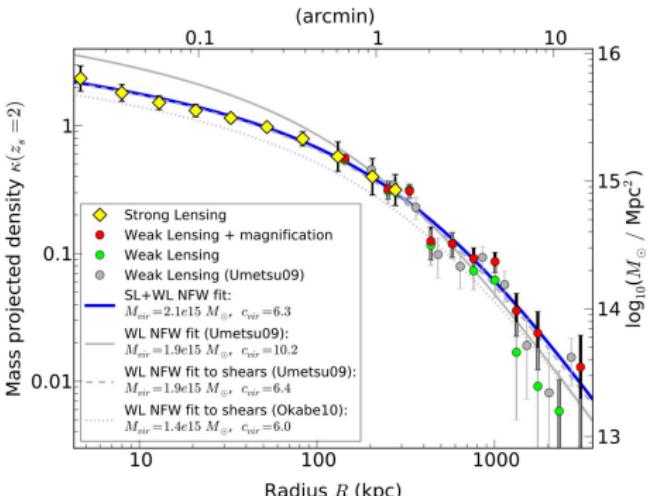


UNIVERSITÄT
HEIDELBERG
ZUKUNFT
SEIT 1386



Abell 2261, CLASH Project

Problems in cosmic structure formation



Coe et al. 2012

Conventional:

- Hydrodynamical equations:

$$\dot{\delta} + \vec{\nabla} \cdot \vec{u} = 0$$

$$\dot{\vec{u}} + 2H\vec{u} = -\vec{\nabla}\phi$$

$$\vec{\nabla}^2\phi = 4\pi G\bar{\rho}\delta$$

- But: **dark matter is no fluid**
- Multiple streams form where shocks would form in a fluid

Conventional:

- Hydrodynamical equations:

$$\dot{\delta} + \vec{\nabla} \cdot \vec{u} = 0$$

$$\dot{\vec{u}} + 2H\vec{u} = -\vec{\nabla}\phi$$

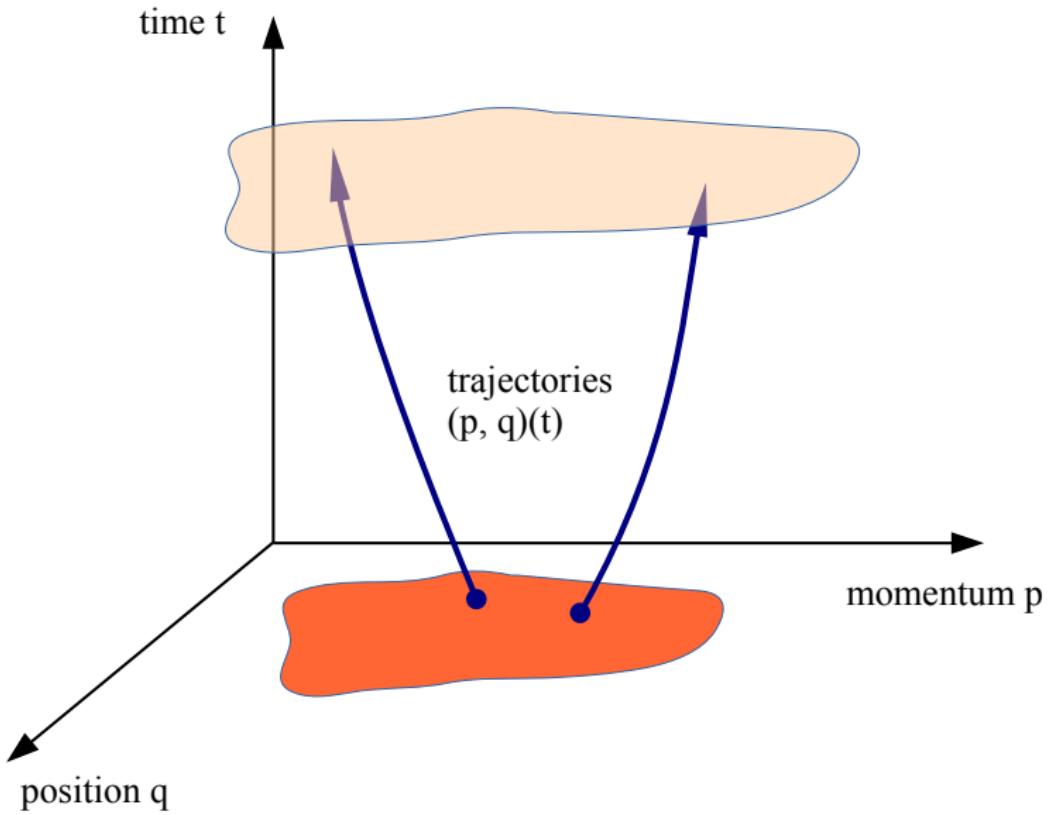
$$\vec{\nabla}^2\phi = 4\pi G\bar{\rho}\delta$$

- But: **dark matter is no fluid**
- Multiple streams form where shocks would form in a fluid

Kinetic Field Theory:

- Non-equilibrium statistics of N classical particle trajectories
- Describe particle ensemble by partition sum (generating functional) Z
- Derive statistical properties by functional derivatives

Phase-space trajectories



Phase-space trajectories

- Classical particles follow Hamiltonian equations of motion,

$$\dot{q} = \frac{\partial \mathcal{H}}{\partial p}, \quad \dot{p} = -\frac{\partial \mathcal{H}}{\partial q}, \quad x = (q, p)$$

- Trajectories are described by (retarded) Green's function $G_R(t, t')$,

$$\bar{x}(t) = \underbrace{G_R(t, 0)x^{(i)}}_{\text{free motion}} - \underbrace{\int_0^t G_R(t, t')\nabla v(t')dt'}_{\text{interaction}}$$



Non-equilibrium, statistical theory for classical degrees of freedom:

$$Z = \int \mathcal{D}\psi P(\psi)$$



Non-equilibrium, statistical theory for classical degrees of freedom:

$$Z = \int \mathcal{D}\psi \int \mathcal{D}\psi^{(i)} P(\psi^{(i)}) P(\psi|\psi^{(i)})$$

Non-equilibrium, statistical theory for classical degrees of freedom:

$$Z = \int \mathcal{D}\psi \int \underbrace{\mathcal{D}\psi^{(i)} P(\psi^{(i)})}_{d\Gamma^{(i)}} P(\psi|\psi^{(i)})$$



Non-equilibrium, statistical theory for classical degrees of freedom:

$$Z[\mathbf{J}] = \int \mathcal{D}\psi \underbrace{\int \mathcal{D}\psi^{(i)} P(\psi^{(i)}) P(\psi|\psi^{(i)})}_{d\Gamma^{(i)}} e^{i\langle \mathbf{J}, \mathbf{x} \rangle}$$

Non-equilibrium, statistical theory for classical degrees of freedom:

$$\begin{aligned} Z[J] &= \int \mathcal{D}\psi \underbrace{\int \mathcal{D}\psi^{(i)} P(\psi^{(i)}) P(\psi|\psi^{(i)})}_{d\Gamma^{(i)}} e^{i\langle J, x \rangle} \\ &= \int dx \int d\Gamma^{(i)} \delta_D [x - \bar{x}(x^{(i)})] e^{i\langle J, x \rangle} \end{aligned}$$

Non-equilibrium, statistical theory for classical degrees of freedom:

$$\begin{aligned} Z[J] &= \int \mathcal{D}\psi \int \underbrace{\mathcal{D}\psi^{(i)} P(\psi^{(i)})}_{d\Gamma^{(i)}} P(\psi|\psi^{(i)}) e^{i\langle J, x \rangle} \\ &= \int dx \int d\Gamma^{(i)} \delta_D [x - \bar{x}(x^{(i)})] e^{i\langle J, x \rangle} \\ &= \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt} \end{aligned}$$

Non-equilibrium, statistical theory for classical degrees of freedom:

$$\begin{aligned} Z[J] &= \int \mathcal{D}\psi \underbrace{\int \mathcal{D}\psi^{(i)} P(\psi^{(i)}) P(\psi|\psi^{(i)}) e^{i\langle J, x \rangle}}_{d\Gamma^{(i)}} \\ &= \int dx \int d\Gamma^{(i)} \delta_D [x - \bar{x}(x^{(i)})] e^{i\langle J, x \rangle} \\ &= \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt} \\ &= e^{i \hat{S}_I} Z_0[J] \end{aligned}$$

Non-equilibrium, statistical theory for classical degrees of freedom:

$$\begin{aligned} Z[J] &= \int \mathcal{D}\psi \int \underbrace{\mathcal{D}\psi^{(i)} P(\psi^{(i)})}_{d\Gamma^{(i)}} P(\psi|\psi^{(i)}) e^{i\langle J, \psi \rangle} \\ &= \int dx \int d\Gamma^{(i)} \delta_D [x - \bar{x}(x^{(i)})] e^{i\langle J, x \rangle} \\ &= \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt} \\ &= e^{i \hat{S}_I} Z_0[J] \end{aligned}$$

$$\langle \rho(1)\rho(2) \rangle = \frac{\delta}{i\delta J(1)} \frac{\delta}{i\delta J(2)} Z[J]$$

MB et al. 2016, 2017, 2019; F. Fabis et al. 2018

Some detail

- Generating functional:

$$Z[J] = e^{i\hat{S}_I} \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt}$$

- Generating functional:

$$Z[J] = e^{i\hat{S}_I} \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt}$$

- Interaction operator:

$$\hat{S}_I = \int d1 \int d2 \hat{\rho}(1) \nabla v(12) \hat{\rho}(2)$$

- Generating functional:

$$Z[J] = e^{i\hat{S}_I} \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt}$$

- Interaction operator:

$$\hat{S}_I = - \int d1 \int d2 \hat{B}(1)v(12)\hat{\rho}(2)$$

- Generating functional:

$$Z[J] = e^{i\hat{S}_I} \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt}$$

- Interaction operator:

$$\hat{S}_I = - \int d1 \int d2 \hat{B}(1)v(12)\hat{\rho}(2)$$

- Perturbation theory: $e^{i\hat{S}_I} = 1 + i\hat{S}_I - \frac{1}{2}\hat{S}_I^2 + \dots$

- Generating functional:

$$Z[J] = e^{i\hat{S}_I} \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt}$$

- Interaction operator:

$$\hat{S}_I = \int d1 \int d2 \hat{\rho}(1) \nabla v(12) \hat{\rho}(2)$$

- Perturbation theory: $e^{i\hat{S}_I} = 1 + i\hat{S}_I - \frac{1}{2}\hat{S}_I^2 + \dots$
- Mean-field approximation: replace \hat{S}_I by $\langle S_I \rangle = \hat{S}_I Z[J] \Big|_{J=0}$

Specialisation to cosmology

- ① Choose initial phase-space measure

$$d\Gamma^{(i)} = P(q, p) d^{3N}q d^{3N}p$$

fully specified by initial power spectrum

- ② Change time coordinate

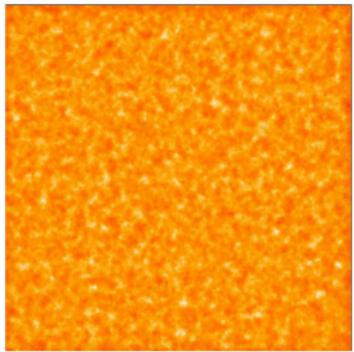
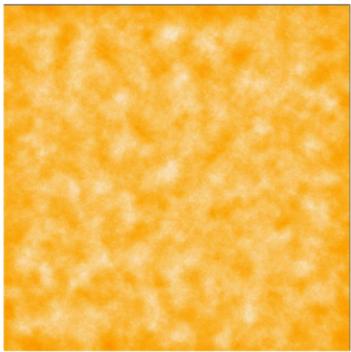
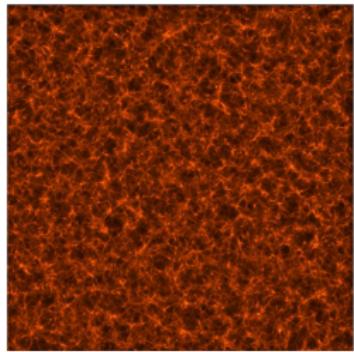
$$t \rightarrow \tau = D_+(t) - D_+(t_i)$$

- ③ Adapt Green's function to expanding universe

Choice of Green's Function



UNIVERSITÄT
HEIDELBERG
ZUKUNFT
SEIT 1386



MB 2015

Interaction in the mean-field approximation



- Instead of perturbative expansion, replace interaction operator by its expectation value:

$$Z[J] = \int d\Gamma^{(i)} e^{i \int \langle J, \bar{x} \rangle dt}$$

with

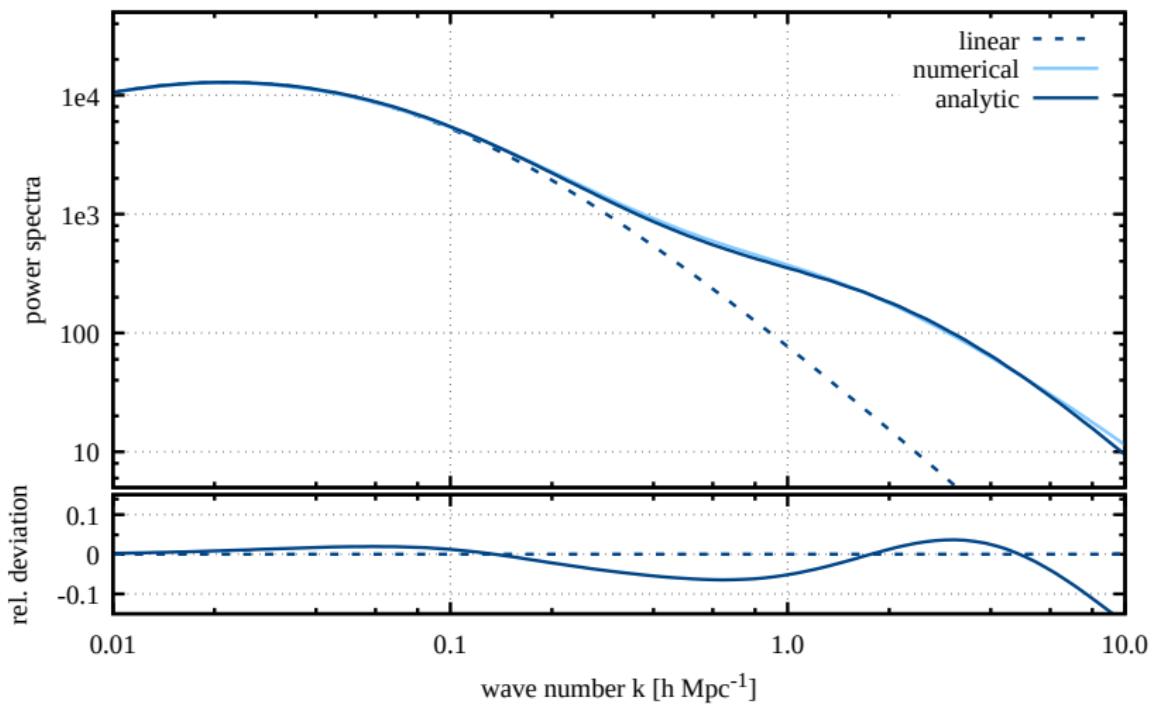
$$\bar{x}(t) = G(t, 0)x^{(i)} - \int_0^t G(t, t') \nabla v(t') dt'$$

- Evaluate ∇v along inertial (unperturbed) trajectories
- Closed expression for non-linear power spectrum:

$$\bar{P}(k, \tau) = e^{Q_D + i \langle S_I \rangle} \int_q \left(e^{\tau^2 k^2 a_{||}(q)} - 1 \right) e^{i \vec{k} \cdot \vec{q}}$$

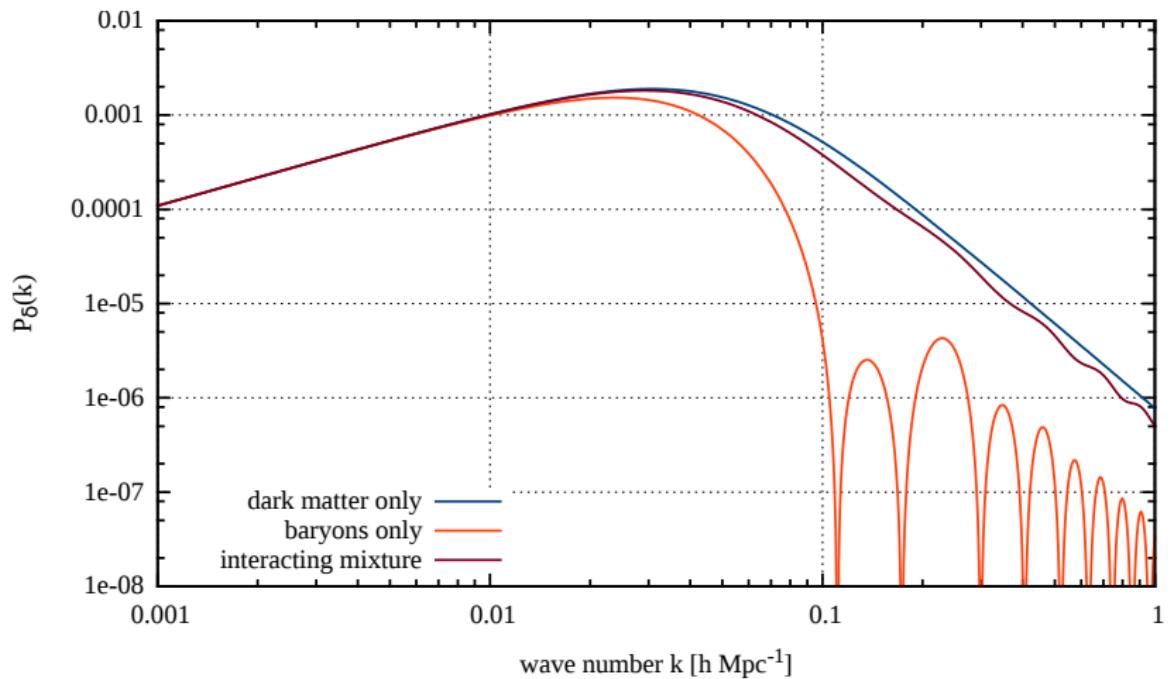
MB et al. 2017; Dombrowski et al. 2018

Density power spectrum

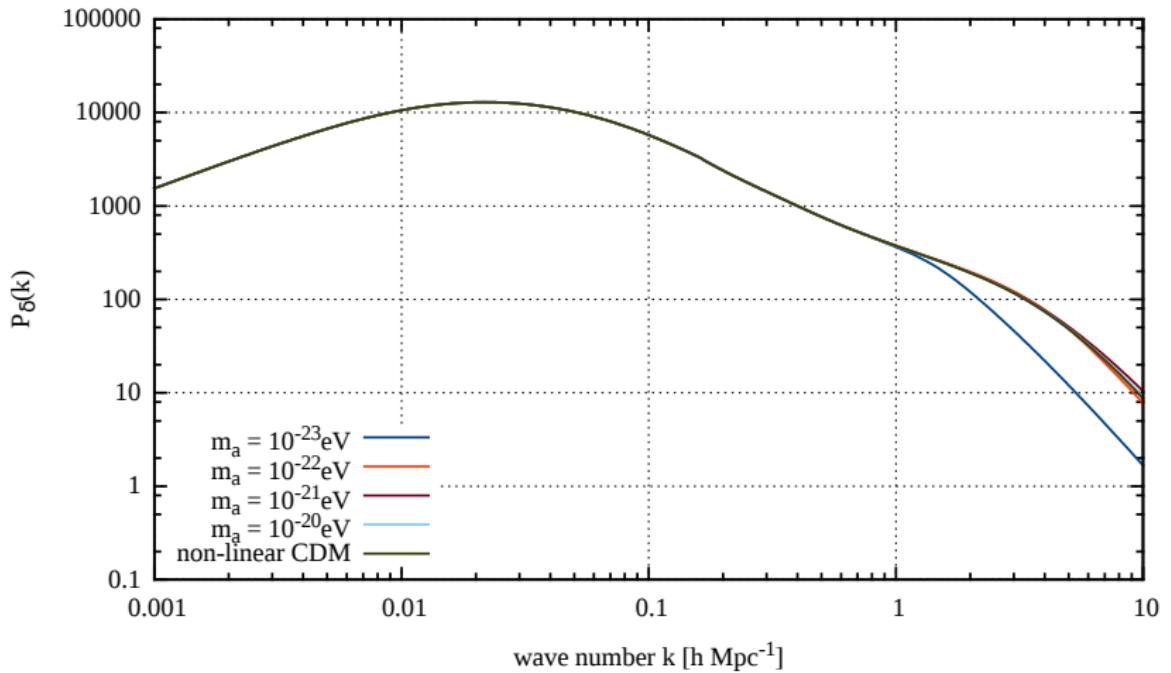


Mean-field approximation (MB et al. 2017, 2019, 2020)

... including gas

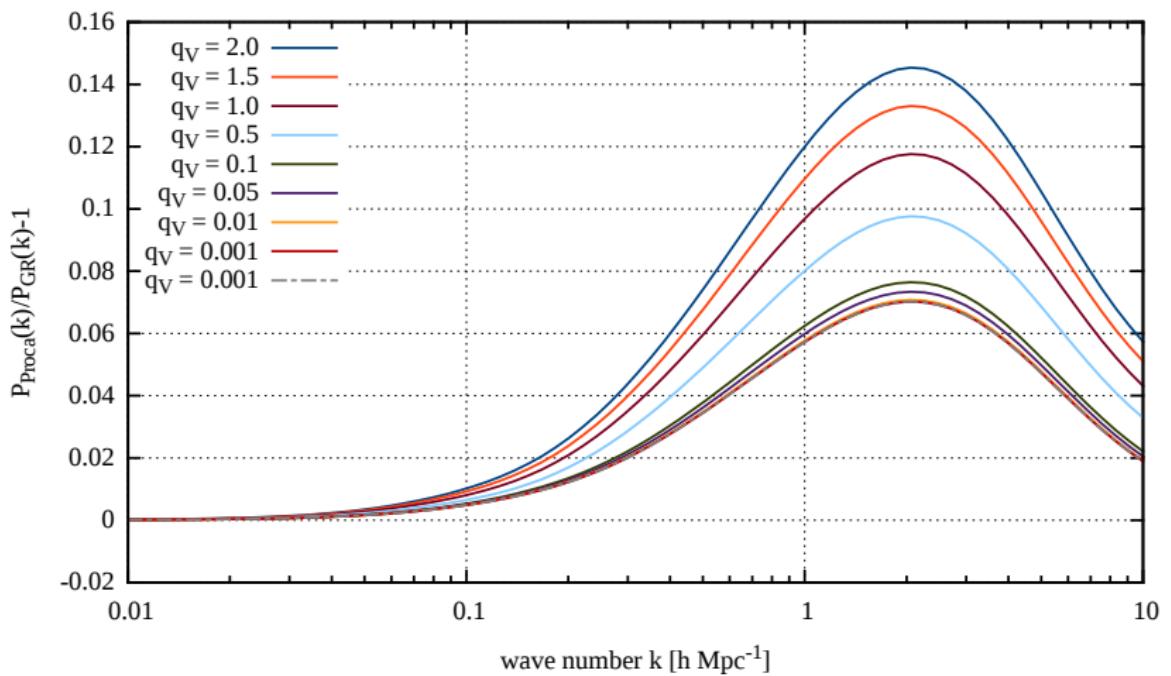


I. Kostyuk, D. Geiss, F. Fabis, R. Lilow, C. Viermann 2016, 2018



C. Littek et al. 2018

... with an alternative field theory of gravity



L. Heisenberg, MB 2018

- ① Non-equilibrium statistical field theory for dark-matter particles set up
- ② Hamiltonian equations of motion, simple Green's function
- ③ Expansion parameter is deviation from unperturbed trajectories
- ④ n -point statistics for collective fields obtained by functional derivatives
- ⑤ Non-perturbative approach with Born approximation reproduces numerical results already quite well
- ⑥ Mixtures of gas and dark matter can be treated in the same way
- ⑦ Generalization to alternative matter models (axions)
- ⑧ Generalization to (some) alternative field theories of gravity straightforward

Recent review: Ann. Phys. 18 (2019) 00446